

Where has all my sand gone?

Hydro-morphodynamics 2D modelling using a discontinuous Galerkin discretisation

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EPSRC
Engineering and Physical Sciences
Research Council



**Imperial College
London**

Overview

Introduction

Building a hydro-morphodynamics 2D model in *Thetis*

Migrating Trench

Meander

Conclusion

Introduction

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Introduction

February 2014 in Dawlish, Devon



Introduction

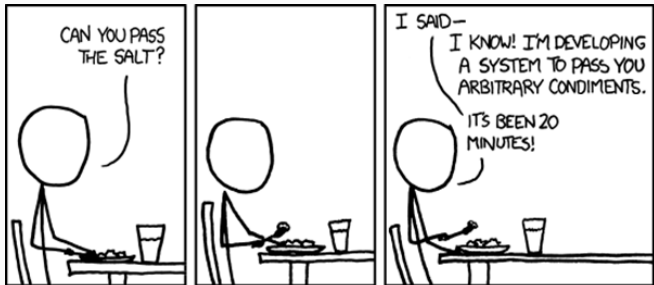
February 2014 in Dawlish, Devon



This cost £35 million to fix and is estimated to have cost the Cornish economy £1.2 billion

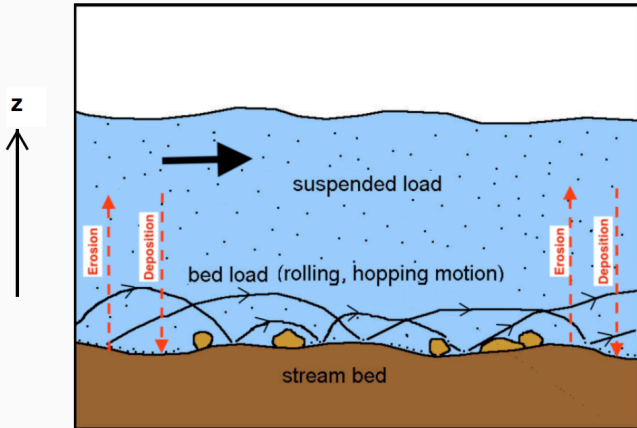
Introduction

Overengineering...



Building a
hydro-morphodynamics 2D
model in *Thetis*

Sediment Transport



Adapted from <http://geologycafe.com/class/chapter11.html>

Basic Model Equations

Depth-averaging from the bed to the water-surface and filtering turbulence:

Hydrodynamics

$$\frac{\partial h}{\partial t} + \frac{\partial}{\partial x}(hU_1) + \frac{\partial}{\partial y}(hU_2) = 0, \quad (1)$$

$$\frac{\partial(hU_j)}{\partial t} + \frac{\partial(hU_jU_1)}{\partial x} + \frac{\partial(hU_jU_2)}{\partial y} = -gh\frac{\partial z_s}{\partial x_j} + \frac{1}{\rho}\frac{\partial(hT_{i1})}{\partial x} + \frac{1}{\rho}\frac{\partial(hT_{i2})}{\partial y} - \frac{\tau_{bi}}{\rho}, \quad (2)$$

Basic Model Equations

Depth-averaging from the bed to the water-surface and filtering turbulence:

Conservation of suspended sediment

$$\frac{\partial}{\partial t}(hC) + \frac{\partial}{\partial x}(hF_{corr}U_1C) + \frac{\partial}{\partial y}(hF_{corr}U_2C) = \frac{\partial}{\partial x} \left[h \left(\epsilon_s \frac{\partial C}{\partial x} \right) \right] + \frac{\partial}{\partial y} \left[h \left(\epsilon_s \frac{\partial C}{\partial y} \right) \right] + E_b - D_b, \quad (1)$$

where z_s is the fluid surface, τ_{bi} the bed shear stress, T_{ij} the depth-averaged stresses, ϵ_s the diffusivity constant and F_{corr} the correction factor.

Calculating the New Bedlevel

Bedlevel (z_b) is governed by the Exner equation

$$\frac{(1 - p')}{m} \frac{dz_b}{dt} + \nabla_h \cdot \mathbf{Q}_b = D_b - E_b, \quad (2)$$

where:

Q_b is the bedload transport given by Meyer-Peter-Müller formula,
 $D_b - E_b$ accounts for effects of suspended sediment flow,
 m is a morphological factor accelerating bedlevel changes.

Adding Physical Effects

Slope Effect

Accounts for gravity which means sediment moves slower uphill than downhill. We impose a magnitude correction:

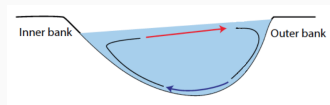
$$Q_{b*} = Q_b \left(1 - \gamma \frac{\partial z_b}{\partial s} \right),$$

and a correction on the flow direction (where δ is the original angle)

$$\tan \alpha = \tan \delta - \tau \frac{\partial z_b}{\partial n}.$$

Secondary Current

Accounts for the helical flow effect in curved channels

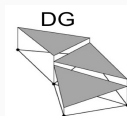


Comparing with Industry Standard Model

Thetis

DG finite element discretisation with

$$P_{1DG} - P_{1DG}$$



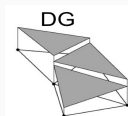
- + Locally mass conservative
- + Well-suited to advection dominated problems
- + Geometrically flexible
- + Allow higher order local approximations

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Telemac-Mascaret

CG finite element discretisation

Method of characteristics
(hydrodynamics advection)

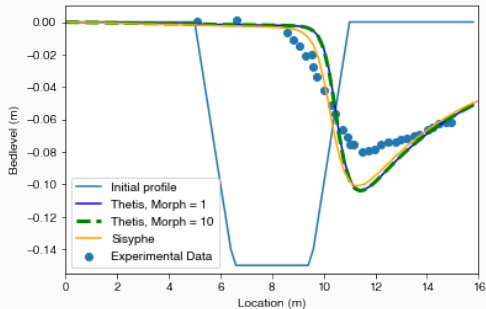
- + Unconditionally stable
- Not mass conservative
- Diffusive for small timesteps

Distributive schemes (sediment transport advection)

- + Mass conservative
- Diffusive for small timesteps
- Courant number limitations to ensure stability

Migrating Trench

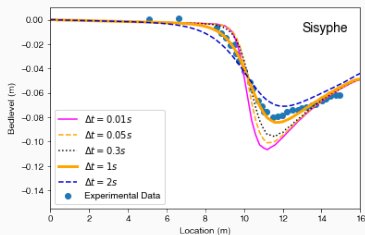
Migrating Trench: Initial Set-up



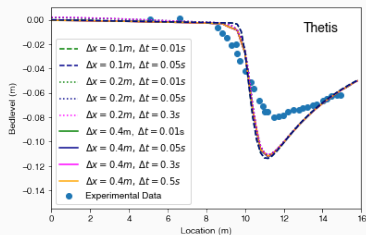
Bedlevel after 15 h for different morphological scale factors comparing experimental data, Sisyphé and Thetis with $\Delta t = 0.05$ s. Experimental data and initial trench profile source: Villaret et al. (2016)

Migrating Trench: Issues with Sisyphé

Varying Δt



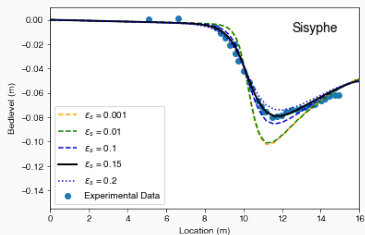
Sisyphé greatly altered by changes to Δt



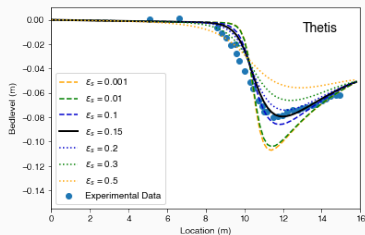
Thetis insensitive to changes in Δt

Migrating Trench: Varying Diffusivity

$$\frac{\partial}{\partial t}(hC) + \frac{\partial}{\partial x}(hF_{corr}U_1C) + \frac{\partial}{\partial y}(hF_{corr}U_2C) =$$
$$\frac{\partial}{\partial x} \left[h \left(\epsilon_s \frac{\partial C}{\partial x} \right) \right] + \frac{\partial}{\partial y} \left[h \left(\epsilon_s \frac{\partial C}{\partial y} \right) \right] + E_b - D_b, \quad (3)$$

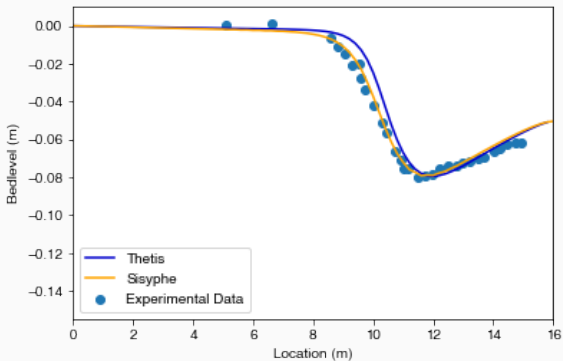


Sensitivity of Sisyphus to ϵ_s



Sensitivity of Thetis to ϵ_s

Migrating Trench: Final Result



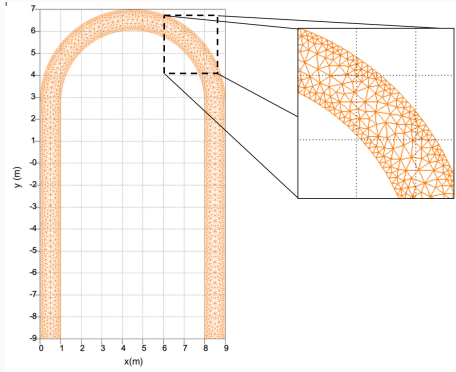
Bedlevel from Thetis and Sisyphus after 15 h

Migrating Trench: Simulation



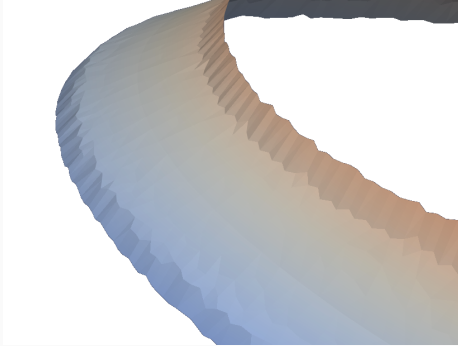
Meander

Meander: Initial Set-up



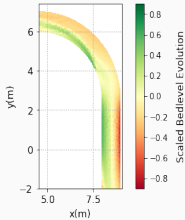
Meander mesh and domain

Meander: Boundary Issue

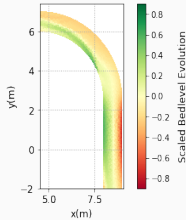


Issue in velocity resolution at boundary resolved by increasing viscosity

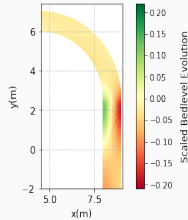
Meander: Physical Effects



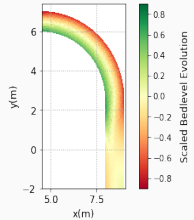
No physical corrections



Only slope effect magnitude

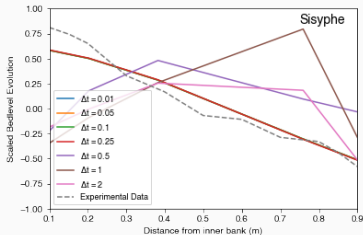


Both slope effect corrections

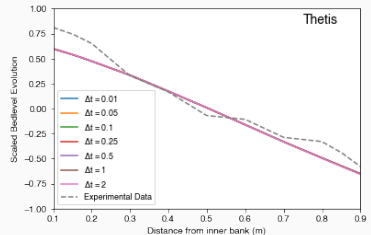


All physical corrections

Meander: Sensitivity to Δt

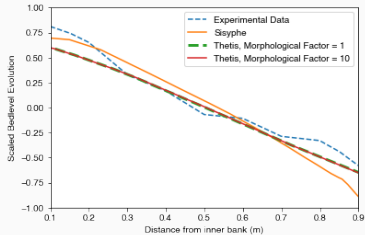


Sisyphus sensitive to changes in Δt

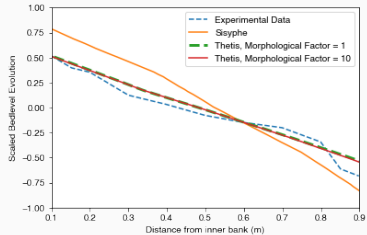


Thetis insensitive to changes in Δt

Meander: Final Result



Cross-section at 90°



Cross-section at 180°

Comparing scaled bedlevel evolution from Thetis, Sisyphé and experimental data

Meander: Simulation



Comparing computational time

	<i>Sisyphé</i>	<i>Thetis</i>	<i>Thetis</i> (morphological scale factor)	<i>Thetis</i> (morphological scale factor, increased Δt)
Migrating Trench	3,427	341,717	39,955	12,422
Meander	980	60,784	10,811	1,212

Comparison of computational time (seconds). For the migrating trench, $\Delta t = 0.05$ s and increased $\Delta t = 0.3$ s; for the meander $\Delta t = 0.1$ s and increased $\Delta t = 10$ s.

Conclusion

Summary

1. Presented the first full morphodynamic model employing a DG based discretisation;

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1. Presented the first full morphodynamic model employing a DG based discretisation;
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3. Validated our model for two different test cases;
4. Shown our model is both accurate and stable, and has key advantages in robustness and accuracy over the state-of-the-art industry standard Syphe whilst still being comparable in computational cost

Key References

-  Kärnä, T., Kramer, S.C., Mitchell, L., Ham, D.A., Piggott, M.D. and Baptista, A.M. (2018), 'Thetis coastal ocean model: discontinuous Galerkin discretization for the threedimensional hydrostatic equations', *Geoscientific Model Development*, **11**, 4359-4382.
-  Tassi, P. and Villaret, C. (2014), *Sisyphé v6.3 User's Manual*, EDF R&D, Chatou, France. Available at:
<http://www.opentelemac.org/downloads/MANUALS/SISYPHE/sisyphé>
-  Villaret, C., Kopmann, R., Wyncoll, D., Riehme, J., Merkel, U. and Naumann, U. (2016), 'First-order uncertainty analysis using Algorithmic Differentiation of morphodynamic models', *Computers & Geosciences*, **90**, 144-151.
-  Villaret, C., Hervouet, J.-M., Kopmann, R., Merkel, U., and Davies, A. G. (2013), 'Morphodynamic modeling using the telemac finite-element system', *Computers & Geosciences*, **53**, 105-113.

Questions?

Using DG:

- Generate a mesh of elements over domain Ω
- Define finite element space on a triangulation (a set of triangles which do not overlap and the union of which is equal to the closure of Ω)
- Derive the weak form of the equation on each triangular element by multiplying the equation by a test function and integrating it by parts on each element and using divergence theorem

Using a discontinuous function space requires the definition of the variables on the element edges thus we use the average and jump operators

$$\{\{X\}\} = \frac{1}{2}(X^+ + X^-), \quad [[\chi]]_n = \chi^+ n^+ + \chi^- n^-, \quad [[X]]_n = X^+ \cdot n^+ + X^- \cdot n^-. \quad (4)$$

For C , we use an upwinding scheme, so, at each edge, C is chosen to be equal to its upstream value with respect to velocity. Therefore

$$\int_{\Omega} \psi \mathbf{u} \cdot \nabla_h C dx = - \int_{\Omega} C \nabla_h \cdot (\mathbf{u} \psi) dx + \int_{\Gamma} C^{\text{up}} [[\psi \mathbf{u}]]_n ds. \quad (5)$$

Weak form of diffusivity term uses Symmetric Interior Penalty Galerkin (SIPG) stabilisation method, as if not discretisation unstable for elliptic operators

$$\begin{aligned} - \int_{\Omega} \psi \nabla_h \cdot (\epsilon_s \nabla_h C) dx &= \int_{\Omega} \epsilon_s (\nabla_h \psi) \cdot (\nabla_h C) dx - \int_{\Gamma} [[\psi]]_n \cdot \{\{\epsilon_s \nabla_h C\}\} ds \\ &\quad - \int_{\Gamma} [[C]]_n \cdot \{\{\epsilon_s \nabla_h \psi\}\} ds + \int_{\Gamma} \sigma \{\{\epsilon_s\}\} [[C]]_n \cdot [[\psi]]_n ds. \end{aligned} \quad (6)$$